

## $L^2$ -ESTIMATES FOR SOME MAXIMAL FUNCTIONS

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SUNTO. Si presentano alcuni risultati sugli operatori massimali associati alle misure superficiali in  $\mathbf{R}^n$ . E. M. Stein [9] ha iniziato lo studio di quest'argomento, dimostrando una disuguaglianza *a priori* per la funzione massimale sferica, mediante l'utilizzo delle funzioni "g". In [3], M. Cowling e G. Mauceri hanno generalizzato il risultato di Stein. Qui si dà una dimostrazione diversa e forse più semplice della stima *a priori*, utilizzando la trasformata di Mellin, come suggerito in un altro lavoro di Cowling e Mauceri [2].

### INTRODUCTION

For a locally integrable function  $f$  on  $\mathbf{R}^n$ , one can define

$$\mathcal{M}f(x) = \sup_{r \in \mathbf{R}^+} \frac{1}{m(B^n)} \int_{B^n} |f(x - ry)| dy, \quad x \in \mathbf{R}^n$$

where  $B^n$  is the unit ball in  $\mathbf{R}^n$  and  $m(B^n)$  is its measure. This is known as the maximal function of  $f$ . One basic property of  $\mathcal{M}f$  is that it satisfies the  $L^p$ -inequality

$$\|\mathcal{M}f\|_p \leq C_p \|f\|_p$$

whenever  $f \in L^p(\mathbf{R}^n)$ ,  $1 < p \leq \infty$ . This was first studied by G. H. Hardy and J. E. Littlewood [4] for  $n = 1$  and, for the general case, by N. Wiener [12] and by J. Marcinkiewicz

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and A. Zygmund [7]. Covering lemmas were used to prove the result. Marcinkiewicz [6] later proved this via an interpolation theorem. For a brief introduction to the theory of maximal functions we refer the reader to Stein [8, Ch. 1].

If  $f$  is smooth and rapidly decreasing, then one can replace  $B^n$  in the definition of  $\mathcal{M}f$  by the unit sphere  $S^{n-1}$  and thus obtain the spherical maximal function of  $f$ . Stein [9] developed the theory and proved that, for  $n \geq 3$ , the *a priori* inequality still holds provided  $f \in L^p(\mathbf{R}^n)$ ,  $\frac{n}{n-1} < p \leq \infty$ . The detailed proof of this can be found in Stein and S. Wainger [10]; it involves the theory of the Fourier transform and an argument using  $g$ -functions.

Cowling and Mauceri [2] reproved Stein's result and gave a new approach to the study of maximal functions. Here the maximal function of  $f$ , denoted now by  $\mathcal{M}_\phi f$ , is defined by

$$\mathcal{M}_\phi f(x) = \sup_{r \in \mathbf{R}^+} |\phi_r * f(x)|, \quad x \in \mathbf{R}^n$$

where  $\phi$  is a distribution on  $\mathbf{R}^n$  and  $\phi_r$  is its dilate; Mellin transform techniques were used to tackle the problems.

A generalisation of Stein's treatment of spherical maximal functions was also found by Cowling and Mauceri [3]. One of their results, namely the  $L^2$ -estimate, asserts that for any smooth function  $f$  on  $\mathbf{R}^n$ , the inequality

$$\| \mathcal{M}_\phi f \|_2 \leq C \| f \|_2$$

holds whenever  $\phi$  is compactly supported in  $\mathbf{R}^n$  and  $|\widehat{\phi}(r\xi)| \leq C(1+r)^{-\alpha}$ ,  $r \in \mathbf{R}^+$ ,  $\xi \in S^{n-1}$ , for some  $\alpha > \frac{1}{2}$ . This was proved via a study of Riesz operators, the theory of fractional differentiation, some properties of Bessel functions, and an argument of  $g$ -functions.

The aim of this note is to prove a similar result for surface measures, using Mellin transform techniques.

Throughout this note we shall use the following notation. The Fourier transform of a function  $f$  on  $\mathbf{R}^n$  is given by

$$\widehat{f}(\xi) = \int_{\mathbf{R}^n} f(x) e^{i\xi x} dx, \quad \xi \in \mathbf{R}^n.$$

$\mathcal{S}(\mathbf{R}^n)$  denotes the space of all smooth and rapidly decreasing functions on  $\mathbf{R}^n$ . By expressions  $C$ ,  $C_k$ ,  $C_{k,l}$  etc. we mean various constants which may vary from line to line. These constants usually depend on  $n$  — the dimension of the ambient space.

Let  $\phi$  be a distribution on  $\mathbf{R}^n$ . For  $r \in \mathbf{R}^+$ , we define the dilate  $\phi_r$  by duality

$$\int_{\mathbf{R}^n} \phi_r(x) f(x) dx = \int_{\mathbf{R}^n} \phi(x) f(rx) dx, \quad f \in \mathcal{S}(\mathbf{R}^n).$$

One may observe that

$$\widehat{\phi}_r(\xi) = \widehat{\phi}(r\xi), \quad r \in \mathbf{R}^+, \quad \xi \in S^{n-1}.$$

Now, for each  $f \in \mathcal{S}(\mathbf{R}^n)$ , define the associated maximal function  $\mathcal{M}_\phi f$  by

$$\mathcal{M}_\phi f(x) = \sup_{r \in \mathbf{R}^+} |\phi_r * f(x)|, \quad x \in \mathbf{R}^n.$$

The focus of our interest is the  $L^2$ -inequality

$$\| \mathcal{M}_\phi f \|_2 \leq C \| f \|_2, \quad f \in \mathcal{S}(\mathbf{R}^n).$$

Using the Mellin transform, we write

$$\widehat{\phi}_r(\xi) = \int_{\mathbf{R}} \psi(u, \xi) r^{iu} du, \quad r \in \mathbf{R}^+,$$

where

$$\psi(u, \xi) = \frac{1}{2\pi} \int_0^\infty \widehat{\phi}_r(\xi) r^{-1-iu} dr, \quad u \in \mathbf{R}.$$

(See [2] for more about the techniques.) For  $s > 0$ , one can observe that

$$\psi(u, s\xi) = s^{iu} \psi(u, \xi), \quad u \in \mathbf{R}, \quad \xi \in \mathbf{R}^n.$$

Moreover, from the theory of the Fourier transform, for all  $f$  in  $\mathcal{S}(\mathbf{R}^n)$ ,

$$(\phi_r * f)^\wedge = \widehat{\phi}_r \widehat{f}$$

whence

$$\begin{aligned} \phi_r * f &= (\widehat{\phi}_r \widehat{f})^\sim \\ &= \left\{ \int_{\mathbf{R}} \psi(u, \cdot) \widehat{f} r^{iu} du \right\}^\sim \\ &= \int_{\mathbf{R}} \{\check{\psi}(u, \cdot) * f\} r^{iu} du. \end{aligned}$$

(Here  $\check{\cdot}$  denotes the inverse Fourier transform.) Hence we find that

$$\mathcal{M}_\phi f(x) \leq \int_{\mathbf{R}} |\check{\psi}(u, \cdot) * f(x)| du, \quad x \in \mathbf{R}^n.$$

Applying Minkowski's inequality and Plancherel's theorem, we obtain

$$\begin{aligned} \|\mathcal{M}_\phi f\|_2 &\leq \int_{\mathbf{R}} \|\check{\psi}(u, \cdot) * f\|_2 du \\ &= \int_{\mathbf{R}} \|\psi(u, \cdot) \widehat{f}\|_2 du \\ &\leq \int_{\mathbf{R}} \|\psi(u, \cdot)\|_\infty \|\widehat{f}\|_2 du \\ &= (2\pi)^{\frac{1}{2}} \|f\|_2 \int_{\mathbf{R}} \|\psi(u, \cdot)\|_\infty du \end{aligned}$$

where

$$\|\psi(u, \cdot)\|_\infty = \sup_{\xi \in S^{n-1}} |\psi(u, \xi)|, \quad u \in \mathbf{R}.$$

Thus, to have the  $L^2$ -estimate, it suffices to verify

$$\int_{\mathbf{R}} \|\psi(u, \cdot)\|_\infty du < \infty.$$

But this would be satisfied when

$$|\psi(u, \xi)| \leq C(1 + |u|)^{-1-\delta}, \quad u \in \mathbf{R}, \xi \in S^{n-1},$$

for some  $\delta > 0$ . With this approach, therefore, we attempt to find conditions on  $\phi$  such that this inequality holds.

We are grateful to Prof. M. Cowling for his suggestions during the preparation and the writing of this paper.

## I

In this part, we are concerned with some maximal operators associated to distributions on  $\mathbf{R}^n$ . Our results here are based on a study of Riesz operators, due to Cowling and Mauceri [3].

For  $a, b$  in  $\mathbf{C}$ , with  $\operatorname{Re}(a) > 0$ , and a distribution  $\phi$  on  $\mathbf{R}^n$ , we define the Riesz operator  $R_{a,b}$  via

$$(R_{a,b}\phi)^\wedge(\xi) = \frac{2}{\Gamma(b)} \int_0^1 s^{a-1} (1-s^2)^{b-1} \widehat{\phi}(s\xi) ds, \quad \xi \in \mathbf{R}^n.$$

This clearly makes sense for  $\operatorname{Re}(b) > 0$ . Moreover, by analytic continuation, it extends to  $\operatorname{Re}(b) > -N$  for any  $N \in \mathbf{Z}^+$  (see [3] for justification).

The following lemma generalises [3, Lemma 1.2].

**Lemma 1.1** Suppose  $\phi$  is a compactly supported distribution on  $\mathbf{R}^n$  and for  $k = 0, 1, 2, \dots$  there exist  $\alpha_k > 0$  such that

$$\left| \frac{\partial^k}{\partial r^k} \widehat{\phi}(r\xi) \right| \leq C_k (1+r)^{-\alpha_k}, \quad r \in \mathbf{R}^+, \xi \in S^{n-1}.$$

Then for  $\operatorname{Re}(a_k) > \alpha_k - k$ ,  $k = 0, 1, 2, \dots$ , and any  $\delta > 0$ , we have

$$\left| \frac{\partial^k}{\partial r^k} (R_{a_k, b}\phi)^\wedge(r\xi) \right| \leq C_k (1+r)^{-\alpha_k + m + \delta}, \quad r \in \mathbf{R}^+, \xi \in S^{n-1},$$

where  $m = \max(0, -\operatorname{Re}(b))$ .

*Proof.* For each  $k = 0, 1, 2, \dots$ , we have, whenever  $r \in \mathbf{R}^+$ ,  $\xi \in S^{n-1}$ ,

$$\begin{aligned} \frac{\partial^k}{\partial r^k} (R_{a_k, b}\phi)^\wedge(r\xi) &= \frac{2}{\Gamma(b)} \int_0^1 s^{a_k-1} (1-s^2)^{b-1} \frac{\partial^k}{\partial r^k} \widehat{\phi}(sr\xi) ds \\ &= \frac{2}{\Gamma(b)} \int_0^1 s^{a_k+k-1} (1-s^2)^{b-1} \frac{\partial^k}{\partial (sr)^k} \widehat{\phi}(sr\xi) ds. \end{aligned}$$

As in the proof of [3, Lemma 1.2], we find that for  $\operatorname{Re}(a_k) > \alpha_k - k$ , and  $\delta > 0$ ,

$$\left| \frac{\partial^k}{\partial r^k} (R_{a_k, b}\phi)^\wedge(r\xi) \right| \leq C_k (1+r)^{-\alpha_k + m + \delta}, \quad r \in \mathbf{R}^+, \xi \in S^{n-1},$$

where  $m = \max(0, -\operatorname{Re}(b))$ .  $\square$

Now we have the following theorem.

**Theorem 1.2** Suppose  $\phi$  is a compactly supported distribution on  $\mathbf{R}^n$  with  $\widehat{\phi}(0) = 0$  such that

$$|\widehat{\phi}(r\xi)| \leq C(1+r)^{-\epsilon}, \quad r \in \mathbf{R}^+, \xi \in S^{n-1}$$

and

$$\left| \frac{\partial}{\partial r} \widehat{\phi}(r\xi) \right| \leq C(1+r)^{-1-\epsilon}, \quad r \in \mathbf{R}^+, \xi \in S^{n-1}$$

for some  $\epsilon > 0$ . Then we have the  $L^2$ -estimate

$$\| \mathcal{M}_\phi f \|_2 \leq C \| f \|_2, \quad f \in \mathcal{S}(\mathbf{R}^n).$$

*Proof.* For each  $\xi$  in  $S^{n-1}$ , write

$$\widehat{\phi}(r\xi) = \int_{\mathbf{R}} \psi(u, \xi) r^{iu} du, \quad r \in \mathbf{R}^+,$$

where

$$\psi(u, \xi) = \frac{1}{2\pi} \int_0^\infty \widehat{\phi}(r\xi) r^{-1-iu} dr, \quad u \in \mathbf{R}.$$

We need to show that  $|\psi(u, \xi)| \leq C(1+|u|)^{-1-\delta}$ ,  $u \in \mathbf{R}$ , for some  $\delta > 0$ .

First, since  $\widehat{\phi}(0) = 0$  and  $\widehat{\phi}$  is differentiable, we have

$$|\widehat{\phi}(r\xi)| \leq Cr, \quad 0 \leq r \leq 1,$$

and thus, whenever  $u \in \mathbf{R}$ , we obtain

$$\begin{aligned} |\psi(u, \xi)| &\leq \int_0^\infty |\widehat{\phi}(r\xi)| r^{-1} dr \\ &\leq \int_0^1 C dr + \int_1^\infty C r^{-1-\epsilon} dr \\ &\leq C. \end{aligned}$$

Next, we invoke the identity

$$\phi = R_{a+2b, -b} R_{a, b} \phi$$

where  $\operatorname{Re}(a) > 0$ ,  $\operatorname{Re}(a+2b) > 0$  (see [3, Lemma 1.3]), to have

$$\frac{\partial}{\partial r} \widehat{\phi}(r\xi) = \frac{2}{\Gamma(-b)} \int_0^1 s^{a+2b-1} (1-s^2)^{-b-1} \frac{\partial}{\partial r} (R_{a, b} \phi) \widehat{\phantom{\phi}}(sr\xi) ds, \quad r \in \mathbf{R}^+.$$

Applying Fubini's theorem, we obtain

$$\begin{aligned}
2\pi i u \psi(u, \xi) &= - \int_0^\infty \widehat{\phi}(r\xi) (-iu) r^{-1-iu} dr \\
&= - \int_0^\infty \widehat{\phi}(r\xi) \frac{\partial}{\partial r} r^{-iu} dr \\
&= \int_0^\infty \frac{\partial}{\partial r} \widehat{\phi}(r\xi) r^{-iu} dr \\
&= \int_0^\infty \frac{2}{\Gamma(-b)} \int_0^1 s^{a+2b-1} (1-s^2)^{-b-1} \frac{\partial}{\partial r} (R_{a,b}\phi)^\wedge(sr\xi) ds r^{-iu} dr \\
&= \frac{2}{\Gamma(-b)} \int_0^1 \int_0^\infty \frac{\partial}{\partial(sr)} (R_{a,b}\phi)^\wedge(sr\xi) (sr)^{-iu} d(sr) s^{a+2b-1+iu} (1-s^2)^{-b-1} ds \\
&= \frac{2}{\Gamma(-b)} \int_0^\infty \frac{\partial}{\partial t} (R_{a,b}\phi)^\wedge(t\xi) t^{-iu} dt \int_0^1 s^{a+2b-1+iu} (1-s^2)^{-b-1} ds \\
&= \frac{\Gamma(\frac{a+2b+iu}{2})}{\Gamma(\frac{a+iu}{2})} \int_0^\infty \frac{\partial}{\partial t} (R_{a,b}\phi)^\wedge(t\xi) t^{-iu} dt.
\end{aligned}$$

For  $\operatorname{Re}(a) > \epsilon$  and  $b = -\frac{\epsilon}{4}$ , we have

$$|u \psi(u, \xi)| \leq \left| \frac{\Gamma(-\frac{\epsilon}{4} + \frac{a+iu}{2})}{\Gamma(\frac{a+iu}{2})} \right| \int_0^\infty \left| \frac{\partial}{\partial t} (R_{a, -\frac{\epsilon}{4}}\phi)^\wedge(t\xi) \right| dt \leq C |u|^{-\frac{\epsilon}{4}}, \quad |u| > 1,$$

since  $|\Gamma(c + id)| \sim C e^{-\frac{\pi}{2}|d|} |d|^{c-\frac{1}{2}}$  as  $|d| \rightarrow \infty$  (see [11, p. 151]) and  $\left| \frac{\partial}{\partial t} (R_{a, -\frac{\epsilon}{4}}\phi)^\wedge(t\xi) \right| \leq C(1+t)^{-1-\frac{\epsilon}{2}}$  for  $t \in \mathbf{R}^+$  (by Lemma 1.1, with  $\delta = \frac{\epsilon}{4}$ ). We therefore find that

$$|\psi(u, \xi)| \leq C |u|^{-1-\frac{\epsilon}{4}}, \quad |u| > 1.$$

Combining this with the previous inequality, we obtain

$$|\psi(u, \xi)| \leq C(1+|u|)^{-1-\frac{\epsilon}{4}}, \quad u \in \mathbf{R}, \quad \xi \in S^{n-1},$$

as desired.  $\square$

*Remark.* The condition  $\widehat{\phi}(0) = 0$  in the theorem can in fact be removed. When  $\phi$  does not satisfy this condition, we can set  $g = \phi - \varphi$ , where  $\varphi \in \mathcal{S}(\mathbf{R}^n)$  with  $\widehat{\varphi}(0) = \widehat{\phi}(0)$ . Thus  $g$  satisfies the hypothesis, and hence the conclusion holds for  $\mathcal{M}_g f$ . But this implies that the conclusion also holds for  $\mathcal{M}_\phi f$  since  $\mathcal{M}_\varphi f$  is known to be dominated by the standard

maximal function of  $f$ . Moreover, the theorem is also true for a distribution  $\phi$  which is not compactly supported but such that  $\widehat{\phi}$  is  $C^1$ .

## II

Here we deal with some maximal operators associated to surface measures on  $\mathbf{R}^n$ .

We first outline the result of Cowling *et al* [1]. Let  $\mathcal{H}$  be a smooth compact convex hypersurface of class  $C^\infty$  in  $\mathbf{R}^n$  whose tangent lines have order of contact at most  $m < \infty$ . Suppose  $\mu$  is a smooth surface measure on  $\mathcal{H}$ . Then, it is shown in [1] that for all  $r \in \mathbf{R}^+$ ,  $\xi \in S^{n-1}$ , we have

$$\widehat{\mu}(r\xi) = F(r, \xi) e^{-irp(\xi)\cdot\xi} + F(r, -\xi) e^{-irp(-\xi)\cdot\xi} + E(r, \xi)$$

where

$$|p(\pm\xi) \cdot \xi| \leq C,$$

$$\left| \frac{\partial^k}{\partial r^k} F(r, \pm\xi) \right| \leq C_k r^{-k} K(r), \quad k = 0, 1, 2, \dots$$

$$\text{with } K(r) = O(r^{-\alpha}) \text{ as } r \rightarrow \infty, \text{ for some } \alpha > 0,$$

$$\text{and } \left| \frac{\partial^k}{\partial r^k} E(r, \xi) \right| \leq C_{k,l} r^{-l}, \quad k, l = 0, 1, 2, \dots$$

By rearranging  $F$  and  $E$  for  $0 \leq r \leq 1$ , we may assume that  $F(0, \pm\xi) = 0$  (or even  $F(r, \pm\xi) = 0$ , for  $0 \leq r < s < 1$ ). If, in addition, we have for all  $r \in \mathbf{R}^+$ ,  $\xi \in S^{n-1}$ ,

$$\left| \frac{\partial^k}{\partial r^k} F(r, \pm\xi) \right| \leq C_k r^{-k-\alpha}, \quad k = 0, 1, 2, \dots,$$

for some  $\alpha > \frac{1}{2}$ , we then say that  $\mu$  is ‘of good decay type’.

Our results are the following.

**Theorem 2.1** Suppose that  $\mathcal{H}$  is a hypersurface in  $\mathbf{R}^n$ , and that  $\mu$  is a surface measure of good decay type on  $\mathcal{H}$ . Then we have the  $L^2$ -estimate

$$\| \mathcal{M}_\mu f \|_2 \leq C \| f \|_2, \quad f \in \mathcal{S}(\mathbf{R}^n).$$



**Theorem 2.2** Let  $\mu$  be a measure on  $\mathbf{R}^n$ ,  $p$  be a bounded function on  $S^{n-1}$ , and  $0 < \epsilon < 3$ . If, for each  $\xi$  in  $S^{n-1}$ , there exists a measure  $\nu = \nu_\xi$  on  $\mathbf{R}$  such that

$$(i) \quad \widehat{\nu}(r) = \widehat{\mu}(r\xi) e^{irp(\xi)\cdot\xi} (1+r^2)^{\frac{1+\epsilon}{4}}, \quad r \in \mathbf{R}^+,$$

$$(ii) \quad \int_{\mathbf{R}} (1+s^2) |d\nu(s)| < C,$$

with  $C$  being independent of  $\xi$ , then we have the  $L^2$ -estimate

$$\| \mathcal{M}_\mu f \|_2 \leq C \| f \|_2, \quad f \in \mathcal{S}(\mathbf{R}^n).$$

We shall prove later that Theorem 2.1 is in fact a special case of Theorem 2.2. Now, Theorem 2.2 follows from the lemma below. We are indebted to Prof. G. Mauceri for suggesting its proof.

**Lemma 2.3** Let  $\mu$  be a measure on  $\mathbf{R}^n$  such that  $\widehat{\mu}(0) = 0$  and  $\nabla \widehat{\mu}(0) = 0$ ,  $p$  be a bounded function on  $S^{n-1}$ , and  $0 < \epsilon < 3$ . Suppose that for each  $\xi$  in  $S^{n-1}$  there exists a measure  $\nu = \nu_\xi$  on  $\mathbf{R}$  such that

$$(i) \quad \widehat{\nu}(r) = \widehat{\mu}(r\xi) e^{irp(\xi)\cdot\xi} (1+r^2)^{\frac{1+\epsilon}{4}}, \quad r \in \mathbf{R}^+,$$

$$(ii) \quad \int_{\mathbf{R}} (1+s^2) |d\nu(s)| < C,$$

with  $C$  being independent of  $\xi$ . Define  $\psi = \psi_\xi$  by

$$\psi(u) = \frac{1}{2\pi} \int_0^\infty \widehat{\mu}(r\xi) r^{-1-iu} dr, \quad u \in \mathbf{R}.$$

Then we have

$$|\psi(u)| \leq C (1+|u|)^{-1-\frac{\epsilon}{2}}, \quad u \in \mathbf{R},$$

with  $C = C(\epsilon, p) \int_{\mathbf{R}} (1+s^2) |d\nu(s)|$ .

*Proof.* We note first that  $\widehat{\nu}$  is  $C^2$  from (ii). We also have  $\widehat{\nu}(0) = 0$  and  $\widehat{\nu}'(0) = 0$ . And, for  $0 \leq r \leq 1$ , we have

$$\begin{aligned}
|\widehat{\nu}(r)| &= |\widehat{\nu}(r) - \widehat{\nu}(0)| \quad (\text{as } \widehat{\nu}(0) = 0) \\
&= \left| r \frac{\partial \widehat{\nu}}{\partial r}(\rho) \right| \quad (\text{for some } 0 < \rho < r) \\
&\leq r \int_{\mathbf{R}} |s| |d\nu(s)| \\
&\leq r \int_{\mathbf{R}} (1 + s^2) |d\nu(s)| \\
&\leq C r.
\end{aligned}$$

Now, from (i), we have

$$\widehat{\mu}(r\xi) = \widehat{\nu}(r) e^{-irq} (1 + r^2)^{-\frac{1+\epsilon}{4}}, \quad r \in \mathbf{R}^+,$$

with  $q = p(\xi) \cdot \xi$ . Thus, whenever  $u \in \mathbf{R}$ ,

$$\begin{aligned}
|\psi(u)| &\leq \int_0^\infty |\widehat{\mu}(r\xi)| r^{-1} dr \\
&\leq \int_0^1 |\widehat{\nu}(r)| r^{-1} dr + \int_1^\infty |\widehat{\nu}(r)| (1 + r^2)^{-\frac{1+\epsilon}{4}} r^{-1} dr \\
&\leq \int_0^1 C dr + \int_1^\infty \|\widehat{\nu}\|_\infty r^{-\frac{3}{2}} dr \\
&\leq C \int_{\mathbf{R}} (1 + s^2) |d\nu(s)| \\
&< \infty.
\end{aligned}$$

From here on we assume  $|u| > 1$ . Take a smooth function  $\varphi$  on  $\mathbf{R}^+$  such that  $\varphi(r) = 0$  if  $0 \leq r \leq \frac{3}{2}$  and  $\varphi(r) = 1$  if  $r \geq 2$ . We then write

$$\begin{aligned}
2\pi \psi(u) &= \int_0^\infty \widehat{\nu}(r) e^{-irq} (1 + r^2)^{-\frac{1+\epsilon}{4}} r^{-1-iu} dr \\
&= \int_0^\infty \widehat{\nu}(r) e^{-irq} \{1 - \varphi(r)\} (1 + r^2)^{-\frac{1+\epsilon}{4}} r^{-1-iu} dr \\
&\quad + \int_0^\infty \widehat{\nu}(r) e^{-irq} \varphi(r) \left\{ (1 + r^2)^{-\frac{1+\epsilon}{4}} - r^{-\frac{1+\epsilon}{2}} \right\} r^{-1-iu} dr \\
&\quad + \int_0^\infty \widehat{\nu}(r) e^{-irq} \varphi(r) r^{-\frac{1+\epsilon}{2}} r^{-1-iu} dr \\
&= \psi_1(u) + \psi_2(u) + \psi_3(u), \quad \text{say.}
\end{aligned}$$

For  $\psi_1(u)$ , we have

$$\begin{aligned} -iu\psi_1(u) &= \int_0^\infty \widehat{\nu}(r) e^{-irq} \{1 - \varphi(r)\} (1+r^2)^{-\frac{1+\epsilon}{4}} \frac{\partial}{\partial r} r^{-iu} dr \\ &= -\int_0^\infty \frac{\partial}{\partial r} \left[ \widehat{\nu}(r) e^{-irq} \{1 - \varphi(r)\} (1+r^2)^{-\frac{1+\epsilon}{4}} \right] r^{-iu} dr, \end{aligned}$$

and similarly

$$-iu(1-iu)\psi_1(u) = \int_0^\infty \frac{\partial^2}{\partial r^2} \left[ \widehat{\nu}(r) e^{-irq} \{1 - \varphi(r)\} (1+r^2)^{-\frac{1+\epsilon}{4}} \right] r^{1-iu} dr.$$

But  $1 - \varphi(r) = 0$  for  $r \geq 2$ , and  $\frac{\partial^2}{\partial r^2} \left[ \widehat{\nu}(r) e^{-irq} \{1 - \varphi(r)\} (1+r^2)^{-\frac{1+\epsilon}{4}} \right]$  is in fact a linear combination of products of derivatives of the four factors, which are all bounded for  $0 \leq r \leq 2$  (in particular, the derivatives of  $\widehat{\nu}$  are dominated by  $\int_{\mathbf{R}} (1+s^2) |d\nu(s)|$ ). Thus we obtain

$$|iu(1-iu)\psi_1(u)| \leq \int_0^2 \left| \frac{\partial^2}{\partial r^2} \left[ \widehat{\nu}(r) e^{-irq} \{1 - \varphi(r)\} (1+r^2)^{-\frac{1+\epsilon}{4}} \right] \right| r dr \leq C,$$

which gives

$$|\psi_1(u)| \leq C |u|^{-2}.$$

For  $\psi_2(u)$ , we also have

$$-iu(1-iu)\psi_2(u) = \int_0^\infty \frac{\partial^2}{\partial r^2} \left[ \widehat{\nu}(r) e^{-irq} \varphi(r) \left\{ (1+r^2)^{-\frac{1+\epsilon}{4}} - r^{-\frac{1+\epsilon}{2}} \right\} \right] r^{1-iu} dr.$$

Put

$$R(r) = (1+r^2)^{-\frac{1+\epsilon}{4}} - r^{-\frac{1+\epsilon}{2}} = r^{-\frac{1+\epsilon}{2}} \left\{ (1+r^{-2})^{-\frac{1+\epsilon}{4}} - 1 \right\}, \quad r \in \mathbf{R}^+.$$

For  $r \geq \frac{3}{2}$ , we can expand  $(1+r^{-2})^{-\frac{1+\epsilon}{4}}$  into its convergent Taylor series to get

$$(1+r^{-2})^{-\frac{1+\epsilon}{4}} = 1 - \frac{1+\epsilon}{4} r^{-2} + \frac{1+\epsilon}{4} \frac{5+\epsilon}{4} \frac{r^{-4}}{2!} - \dots.$$

Hence

$$R(r) = -\frac{1+\epsilon}{4} r^{-\frac{5+\epsilon}{2}} + \text{terms of a lower power of } r, \quad r \geq \frac{3}{2}.$$

We therefore find that  $|R(r)|$ ,  $|\frac{\partial}{\partial r}R(r)|$  and  $|\frac{\partial^2}{\partial r^2}R(r)| \leq C r^{-\frac{5+\epsilon}{2}}$  for  $r \geq \frac{3}{2}$ . And thus, since the derivatives of  $\widehat{\nu}$ ,  $e^{-irq}$  and  $\varphi$  are bounded, we obtain

$$|iu(1-iu)\psi_2(u)| \leq \int_{\frac{3}{2}}^{\infty} C r^{-\frac{5+\epsilon}{2}} dr \leq C,$$

which gives

$$|\psi_2(u)| \leq C|u|^{-2}.$$

For  $\psi_3(u)$ , we write

$$\begin{aligned} \psi_3(u) &= \int_0^{\infty} \widehat{\nu}(r) e^{-irq} \varphi(r) r^{-\frac{3+\epsilon}{2}-iu} dr \\ &= \int_0^{\infty} \widehat{\nu}(r) e^{-irq} \varphi(r) r^{z-1} dr \quad (\text{with } z = -\frac{1+\epsilon}{2} - iu) \\ &= \int_{\mathbf{R}} \int_0^{\infty} \varphi(r) r^{z-1} e^{-ir(s+q)} dr d\nu(s) \quad (\text{as } \varphi(r) = 0 \text{ for } 0 \leq r \leq \frac{3}{2}). \end{aligned}$$

Then, for fixed  $s \in \mathbf{R}$ , consider

$$I(z) = \int_0^{\infty} \varphi(r) r^{z-1} e^{-ir(s+q)} dr, \quad z \in \mathbf{C}.$$

We note that  $I(z)$  continues analytically into  $\text{Re}(z) \leq 1$ . For  $0 < \text{Re}(z) < 1$ , we write

$$\begin{aligned} I(z) &= \int_0^{\infty} \{\varphi(r) - 1\} r^{z-1} e^{-ir(s+q)} dr + \int_0^{\infty} r^{z-1} e^{-ir(s+q)} dr \\ &= I_4(z) + I_5(z), \quad \text{say.} \end{aligned}$$

Corresponding to  $I_4(z)$ , we have

$$\psi_4(z) = \int_{\mathbf{R}} I_4(z) d\nu(s) = \int_0^{\infty} \widehat{\nu}(r) e^{-irq} \{\varphi(r) - 1\} r^{z-1} dr,$$

which continues analytically into  $-2 < \text{Re}(z) < 1$ , since  $\widehat{\nu}(0) = 0$  and  $\widehat{\nu}'(0) = 0$ . Hence

$$\begin{aligned} \psi_4(u) &= \int_0^{\infty} \widehat{\nu}(r) e^{-irq} \{\varphi(r) - 1\} r^{-\frac{3+\epsilon}{2}-iu} dr \\ &= \int_0^{\infty} \left[ \widehat{\nu}(r) e^{-irq} \{\varphi(r) - 1\} r^{-\frac{1+\epsilon}{2}} \right] r^{-1-iu} dr. \end{aligned}$$

Integrating by parts twice, we obtain

$$-iu(1-iu)\psi_4(u) = \int_0^{\infty} \frac{\partial^2}{\partial r^2} \left[ \widehat{\nu}(r) e^{-irq} \{\varphi(r) - 1\} r^{-\frac{1+\epsilon}{2}} \right] r^{1-iu} dr.$$

But  $\varphi(r) - 1 = 0$  for  $r \geq 2$ , and  $|\frac{\partial^2}{\partial r^2} [\widehat{\nu}(r) e^{-irq} \{\varphi(r) - 1\} r^{-\frac{1+\epsilon}{2}}]| \leq C r^{-\frac{1+\epsilon}{2}}$  for  $0 \leq r \leq 2$  since again the derivatives of  $\widehat{\nu}$ ,  $e^{-irq}$  and  $\varphi$  are bounded. Hence we find that

$$|iu(1 - iu) \psi_4(u)| \leq \int_0^2 C r^{\frac{1-\epsilon}{2}} dr < \infty \quad (\text{as } 0 < \epsilon < 3),$$

whence

$$|\psi_4(u)| \leq C |u|^{-2}.$$

It remains to estimate  $\psi_5(u) = \int_{\mathbf{R}} I_5(z) d\nu(s)$ , where  $z = -\frac{1+\epsilon}{2} - iu$ . Let  $h$  be the function on  $\mathbf{C}$  defined by

$$h(w) = w^{z-1} e^{-wt}, \quad z = -\frac{1+\epsilon}{2} - iu, \quad t \in \mathbf{R}.$$

By integrating  $h$  around the contour  $\gamma_1$  (see Fig. 1) for  $t > 0$ , or around the contour  $\gamma_2$  (see Fig. 2) for  $t < 0$ , one can observe that, as  $R \rightarrow \infty$ ,

$$\int_0^\infty r^{z-1} e^{-irt} dr = (it)^{-z} \Gamma(z), \quad t \neq 0.$$

Fig. 1 The contour  $\gamma_1$

Fig. 2 The contour  $\gamma_2$

Hence

$$\begin{aligned} |\psi_5(u)| &= \left| \int_{\mathbf{R}} I_5(z) d\nu(s) \right| && (\text{with } z = -\frac{1+\epsilon}{2} - iu) \\ &= \left| \int_{\mathbf{R}} \int_0^\infty r^{z-1} e^{-ir(s+q)} dr d\nu(s) \right| \\ &= \left| \int_{\mathbf{R}} \{i(s+q)\}^{-z} \Gamma(z) d\nu(s) \right| \\ &\leq e^{-\frac{\pi}{2}u} |\Gamma(-\frac{1+\epsilon}{2} - iu)| \int_{\mathbf{R}} |s+q|^{\frac{1+\epsilon}{2}} |d\nu(s)| \\ &\leq C |u|^{-1-\frac{\epsilon}{2}} \int_{\mathbf{R}} |s+q|^{\frac{1+\epsilon}{2}} |d\nu(s)| && (\text{as } |u| > 1) \\ &\leq C |u|^{-1-\frac{\epsilon}{2}} \int_{\mathbf{R}} \{1 + (s+q)^2\} |d\nu(s)| && (\text{as } 0 < \epsilon < 3) \\ &\leq C |u|^{-1-\frac{\epsilon}{2}} (1+q^2) \int_{\mathbf{R}} (1+s^2) |d\nu(s)| \\ &\leq C |u|^{-1-\frac{\epsilon}{2}}. \end{aligned}$$

Combining this with the previous results, we obtain

$$|\psi(u)| \leq C |u|^{-1-\frac{\epsilon}{2}}, \quad |u| > 1,$$

with  $C = C(\epsilon, p) \int_{\mathbf{R}} (1 + s^2) |d\nu(s)|$ . This completes the proof.  $\square$

To prove Theorem 2.1, we adapt a result of Hörmander [5, pp. 121-122].

**Proposition 2.4** Suppose that  $\Phi$  is a smooth function on  $\mathbf{R}$ , and that for  $k = 0, 1, 2, \dots$

$$\left| \frac{\partial^k}{\partial r^k} \Phi(r) \right| \leq C_k (1 + |r|)^{-k-\epsilon}, \quad r \in \mathbf{R},$$

for some  $\epsilon > 0$ . It then follows that  $\check{\Phi} \in L^1(\mathbf{R})$ . In general,  $s^k \check{\Phi}(s) \in L^1(\mathbf{R})$ , for  $k = 0, 1, 2, \dots$ .

*Proof.* First of all (see [5, p. 121]), there exists a function  $\varphi \in C_0^\infty(\mathbf{R})$  supported in  $\{r : \frac{1}{2} < |r| < 2\}$  such that

$$\sum_{j=-\infty}^{\infty} \varphi(2^{-j}r) = 1, \quad r \neq 0.$$

For this  $\varphi$ , we have

$$\sum_{j=1}^{\infty} \varphi(2^{-j}r) = 0, \quad 0 < |r| \leq \frac{1}{2}$$

and

$$\sum_{j=1}^{\infty} \varphi(2^{-j}r) = 1, \quad |r| > 2.$$

Let us put

$$\varphi_0(r) = \begin{cases} 1 - \sum_{j=1}^{\infty} \varphi(2^{-j}r), & \text{if } r \neq 0, \\ 1, & \text{if } r = 0. \end{cases}$$

It is clear that  $\varphi_0 \in C_0^\infty(\mathbf{R})$  and that

$$\varphi_0(r) + \sum_{j=1}^{\infty} \varphi(2^{-j}r) = 1, \quad r \in \mathbf{R}.$$

This enables us to decompose  $\Phi$  into

$$\Phi = \Phi_0 + \sum_{j=1}^{\infty} \Phi_j$$

where  $\Phi_0 = \varphi_0 \Phi$  and  $\Phi_j(r) = \varphi(2^{-j}r) \Phi(r)$ ,  $r \in \mathbf{R}$  ( $j = 0, 1, 2, \dots$ ).

We see that  $\Phi_0 = \varphi_0 \Phi$  is smooth and compactly supported on  $\mathbf{R}$ . Thus for all  $s \in \mathbf{R}$

$$|\check{\Phi}_0(s)| \leq \frac{1}{2\pi} \int_{\mathbf{R}} |\Phi_0(r)| dr \leq C$$

and

$$|s^2 \check{\Phi}_0(s)| \leq \frac{1}{2\pi} \int_{\mathbf{R}} \left| \frac{\partial^2}{\partial r^2} \Phi_0(r) \right| dr \leq C.$$

Hence

$$|\check{\Phi}_0(s)| \leq C(1 + |s|)^{-2}, \quad s \in \mathbf{R},$$

yielding

$$\int_{\mathbf{R}} |\check{\Phi}_0(s)| ds < \infty.$$

Now, for  $j = 1, 2, 3, \dots$ , we have

$$\begin{aligned} \int_{\mathbf{R}} |\Phi_j(r)|^2 dr &= \int_{\mathbf{R}} |\varphi(2^{-j}r) \Phi(r)|^2 dr \\ &\leq C \int_{2^{j-1} \leq |r| \leq 2^{j+1}} |\Phi(r)|^2 dr \\ &\leq C \int_{2^{j-1} \leq |r| \leq 2^{j+1}} |r|^{-2\epsilon} dr \\ &\leq C 2^{(1-2\epsilon)j}, \end{aligned}$$

and similarly

$$\begin{aligned} &\int_{\mathbf{R}} \left| 2^j \frac{\partial}{\partial r} \Phi_j(r) \right|^2 dr \\ &= \int_{\mathbf{R}} \left| \frac{\partial \varphi}{\partial r}(2^{-j}r) \Phi(r) + 2^j \varphi(2^{-j}r) \frac{\partial}{\partial r} \Phi(r) \right|^2 dr \\ &\leq C \int_{2^{j-1} \leq |r| \leq 2^{j+1}} \left\{ |\Phi(r)|^2 + 2^j |\Phi(r)| \left| \frac{\partial}{\partial r} \Phi(r) \right| + 2^{2j} \left| \frac{\partial}{\partial r} \Phi(r) \right|^2 \right\} dr \\ &\leq C \int_{2^{j-1} \leq |r| \leq 2^{j+1}} \left\{ |r|^{-2\epsilon} + 2^j |r|^{-1-2\epsilon} + 2^{2j} |r|^{-2-2\epsilon} \right\} dr \\ &\leq C 2^{(1-2\epsilon)j}. \end{aligned}$$

Then, by Plancherel's theorem,

$$\int_{\mathbf{R}} |\check{\Phi}_j(s)|^2 ds = \int_{\mathbf{R}} |\Phi_j(r)|^2 dr \leq C 2^{(1-2\epsilon)j}$$

and

$$\int_{\mathbf{R}} |2^j s \check{\Phi}_j(s)|^2 ds = \int_{\mathbf{R}} |2^j \frac{\partial}{\partial r} \Phi_j(r)|^2 dr \leq C 2^{(1-2\epsilon)j}.$$

These give

$$\int_{\mathbf{R}} 2^{-j} (1 + 2^{2j} s^2) |\check{\Phi}_j(s)|^2 ds \leq C 2^{-2\epsilon j}.$$

Applying the Cauchy-Schwartz inequality, we obtain

$$\begin{aligned} \int_{\mathbf{R}} |\check{\Phi}_j(s)| ds &\leq \left\{ \int_{\mathbf{R}} 2^{-j} (1 + 2^{2j} s^2) |\check{\Phi}_j(s)|^2 ds \right\}^{\frac{1}{2}} \left\{ \int_{\mathbf{R}} 2^j (1 + 2^{2j} s^2)^{-1} ds \right\}^{\frac{1}{2}} \\ &\leq C 2^{-\epsilon j} \left\{ \int_{|s| \leq 2^{-j}} 2^j ds + \int_{2^{-j} \leq |s| \leq 1} ds + \int_{|s| \geq 1} 2^{-j} s^{-2} ds \right\}^{\frac{1}{2}} \\ &\leq C 2^{-\epsilon j}. \end{aligned}$$

Since  $\check{\Phi} = \check{\Phi}_0 + \sum_{j=1}^{\infty} \check{\Phi}_j$ , we therefore have

$$\|\check{\Phi}\|_1 \leq \|\check{\Phi}_0\|_1 + \sum_{j=1}^{\infty} \|\check{\Phi}_j\|_1 \leq C(\epsilon) < \infty,$$

meaning that  $\check{\Phi} \in L^1(\mathbf{R})$ .

In general, the proof also applies to  $\frac{\partial^k}{\partial r^k} \Phi(r)$ , and so we have  $s^k \check{\Phi}(s) \in L^1(\mathbf{R})$ , for all  $k = 0, 1, 2, \dots$ .  $\square$

We now come to the proof of Theorem 2.1.

*Proof* (of Theorem 2.1). For  $r \in \mathbf{R}^+$ ,  $\xi \in S^{n-1}$ , we have

$$\widehat{\mu}(r\xi) = F(r, \xi) e^{-irp(\xi)\cdot\xi} + F(r, -\xi) e^{-irp(-\xi)\cdot\xi} + E(r, \xi),$$

with  $p$ ,  $F$ , and  $E$  as prescribed. We assume here that  $F(r, \pm\xi) = 0$ , for  $0 \leq r < s < 1$ , to have  $\frac{\partial^k}{\partial r^k} F(r, \xi)|_{r=0} = 0$ , for all  $k = 0, 1, 2, \dots$ .



For  $u \in \mathbf{R}$ ,  $\xi \in S^{n-1}$ , consider

$$\begin{aligned}
\psi(u, \xi) &= \frac{1}{2\pi} \int_0^\infty \widehat{\mu}(r\xi) r^{-1-iu} dr \\
&= \frac{1}{2\pi} \int_0^\infty F(r, \xi) e^{-irp(\xi)\cdot\xi} r^{-1-iu} dr \\
&\quad + \frac{1}{2\pi} \int_0^\infty F(r, -\xi) e^{-irp(-\xi)\cdot\xi} r^{-1-iu} dr \\
&\quad + \frac{1}{2\pi} \int_0^\infty E(r, \xi) r^{-1-iu} dr \\
&= \psi_1(u, \xi) + \psi_2(u, \xi) + \psi_3(u, \xi), \quad \text{say.}
\end{aligned}$$

Let us fix  $\xi \in S^{n-1}$  hereafter. By Theorem 1.2, we have

$$|\psi_3(u, \xi)| \leq C (1 + |u|)^{-1-\delta}, \quad u \in \mathbf{R},$$

for some  $\delta > 0$ , assuming that  $E(0, \xi) = 0$ .

It then remains to tackle  $\psi_1$  and  $\psi_2$ . But since they are similar, it suffices to work with one of them,  $\psi_1$  say. We put

$$\widehat{\mu}_1(r\xi) = F(r, \xi) e^{-irp(\xi)\cdot\xi},$$

and define  $\widehat{\nu}$  on  $\mathbf{R}$  by

$$\widehat{\nu}(r) = \begin{cases} \widehat{\mu}_1(r\xi) e^{irp(\xi)\cdot\xi} (1+r^2)^{\frac{1+\epsilon}{4}}, & \text{if } r \geq 0, \\ \widehat{\mu}_1((-r)(-\xi)) e^{irp(-\xi)\cdot\xi} (1+r^2)^{\frac{1+\epsilon}{4}}, & \text{if } r < 0, \end{cases}$$

for some  $0 < \epsilon < \min(3, \alpha - \frac{1}{2})$ . Thus

$$\widehat{\nu}(r) = \begin{cases} F(r, \xi) (1+r^2)^{\frac{1+\epsilon}{4}}, & \text{if } r \geq 0, \\ F(-r, -\xi) (1+r^2)^{\frac{1+\epsilon}{4}}, & \text{if } r < 0. \end{cases}$$

Writing

$$\widehat{\nu}(r) = \begin{cases} F(r, \xi) r^{\frac{1+\epsilon}{2}} (1+r^{-2})^{\frac{1+\epsilon}{4}}, & \text{if } r \geq 0, \\ F(-r, -\xi) (-r)^{\frac{1+\epsilon}{2}} (1+r^{-2})^{\frac{1+\epsilon}{4}}, & \text{if } r < 0, \end{cases}$$

and applying Leibnitz's formula, we obtain that for  $k = 0, 1, 2, \dots$ ,

$$\left| \frac{\partial^k}{\partial r^k} \widehat{\nu}(r) \right| \leq C, \quad |r| \leq 1$$

and

$$\left| \frac{\partial^k}{\partial r^k} \widehat{\nu}(r) \right| \leq C |r|^{-k-\frac{\epsilon}{2}}, \quad |r| > 1.$$

Hence, for all  $k = 0, 1, 2, \dots$ , we have

$$\left| \frac{\partial^k}{\partial r^k} \widehat{\nu}(r) \right| \leq C (1 + |r|)^{-k - \frac{\epsilon}{2}}, \quad r \in \mathbf{R}.$$

It follows from Proposition 2.4 that  $\nu = (\widehat{\nu})^\sim \in L^1(\mathbf{R})$  (which assures us that  $\nu$  defines a measure on  $\mathbf{R}$ ), and further  $s^2 \nu(s) \in L^1(\mathbf{R})$ . So we find that there exists a measure  $\nu$  on  $\mathbf{R}$  which satisfies the hypothesis of Lemma 2.3. The proof is therefore complete.  $\square$

*Remark.* The result extends to surface measures of class  $C^K$  for some finite  $K$ ;  $K$  is sufficiently large so that the hypothesis

$$\left| \frac{\partial^k}{\partial r^k} F(r, \pm \xi) \right| \leq C_k r^{-k - \alpha}, \quad r \in \mathbf{R}^+, \quad \xi \in S^{n-1},$$

is satisfied for  $k = 0, 1, 2, 3$  and 4.

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